

Fig. 4 Activity of constraints for mild steel beam column.

under suppression of local buckling and yielding the optimum mild steel tube has active local and yield constraints with inherent stiffness considerably in excess of a 1% deflection limitation. The discontinuous behavior of  $t/L$  and  $D/L$  where  $P/L^2 < 0.35$  for the case  $Q/L^2 = 0$  in Fig. 3 accompanied by continuity in weight merit (Fig. 2) shows the existence of multiple optima for the simple column beyond the load index for which the yield constraint becomes active. The corresponding inactivity of either the local or Euler modes or a combination gives rise to a family of cross sections each having the property of least weight. Although no unique pair ( $t$ ,  $D$ ) can be found from the programmed solution, their product is computable so as to achieve the yielding area, and then either the local or Euler constraints can be set equal to yield stress to effect a unique sizing.

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## Floating Piecewise Linear Approximation of a Nonlinear Constitutive Equation

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#### Introduction

ONE of the main difficulties arising in large deflection analysis of structures has its origin in the nonlinearity of the stress-strain relations. This difficulty is usually overcome by piecewise linearization of the constitutive equation. A method for construction of the "best" piecewise linear approximation is proposed. It will be shown, through an examination of a simple model, that the suggested method yields a sequence of upper bounds to the load-deflection behavior. The maximum load carrying capacities for the given model obtained by the approximate and the exact methods are the same.

#### Exact Solution to a Simple Model

Consider the simple model shown in Fig. 1. This model is taken from Ref. 1 where an approximate analysis has been given. The model consists of a rigid bar  $AB$  and a spring  $BC$ . The stress strain relation of the spring is given by

$$\theta - \theta_0 = (\sigma/E) + K(\sigma/E)^n \quad (1)$$

where  $\sigma$  is the stress in the spring and  $\theta_0$  is an initial imperfection. Equation (1) is of the type given in Ref. 2. The static equilibrium requirement, for small value of  $\theta$  ( $t\theta \approx 0$ ) is

$$A\sigma = P\theta \quad (2)$$

where  $A$  is the area of the spring and  $P$  the external applied load. By substitution of Eq. (2) into Eq. (1), the exact load-deflection behavior is obtained

$$\theta - \theta_0 = (P\theta/EA) + K(P\theta/EA)^n \quad (3)$$

The behavior of the perfect system is obtained from Eq. (3) by taking  $\theta_0 = 0$ . The resulting nonlinear equation has two branches

$$\theta = 0 \quad \text{and} \quad (P/EA)[1 + K(P\theta/EA)^{n-1}] = 1 \quad (4)$$

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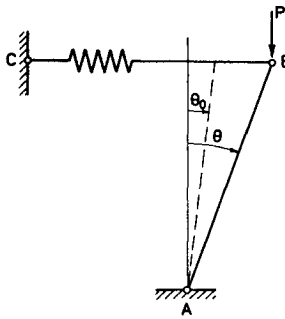


Fig. 1 A simple model.

with the bifurcation point

$$\theta_{cr} = 0 \quad P_{cr}/EA = 1 \quad (5)$$

The imperfect system Eq. (3) has no branching point and its instability is obtained at the maximum possible load  $P^*$  (at which  $dP/d\theta = 0$ ). Differentiating Eq. (3) yields

$$1 = \frac{P^*}{P_{cr}} \left[ 1 + nK \left( \frac{P^*}{P_{cr}} \theta^* \right)^{n-1} \right] \quad (6)$$

and also

$$\theta^* - \theta_o = \frac{P^*}{P_{cr}} \theta^* + K \left( \frac{P^*}{P_{cr}} \theta^* \right)^n \quad (7)$$

where  $\theta^*$  is the angle corresponding to  $P^*$ . Fortunately Eqs. (6) and (7) admit a simple solution

$$\frac{P^*}{P_{cr}} = 1 / \left[ 1 + nK^{1/n} \left( \frac{\theta_o}{n-1} \right)^{(n-1)/n} \right] \quad (8)$$

$$\theta^* = \frac{n\theta_o}{n-1} + \left[ \frac{\theta_o}{(n-1)K} \right]^{1/n} \quad (9)$$

Typical results for the perfect and the imperfect system are shown in Figs. 2 and 3.

#### The Lowest Modulus Concept

Before proceeding, the concept of the lowest modulus will be introduced. This modulus will be used later in the approximate analysis of the model.

Consider again the Ramberg Osgood stress strain relation

$$\varepsilon = (\sigma/E) + K(\sigma/E)^n \quad (10)$$

and suppose that Eq. (10) is replaced by the two tangents OA and ABC (Fig. 4). The slope of the first tangent is the regular

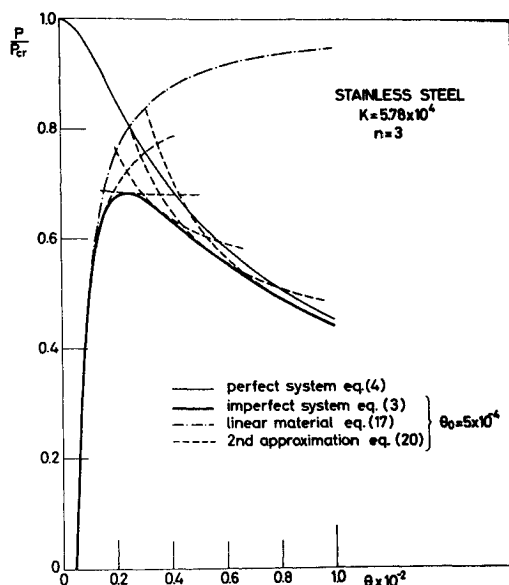


Fig. 2 Load deflection behavior, stainless steel.

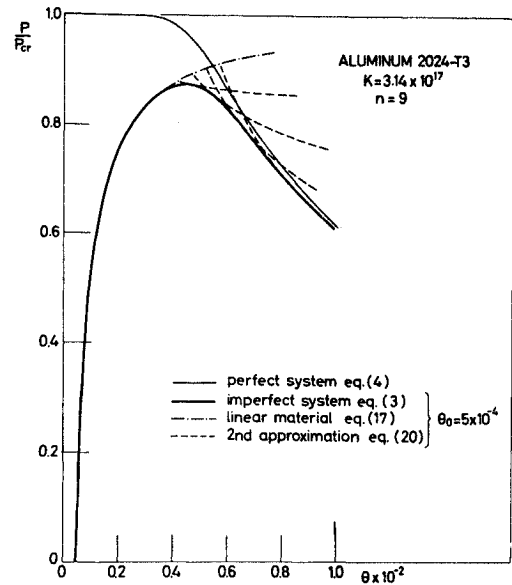


Fig. 3 Load deflection behavior, 2024-T3 Aluminum.

modulus of elasticity  $E$ . Let the slope of the second tangent be  $E_L$ . It is understood, at least intuitively, that the two tangents represent a stiffer material than the one given by Eq. (10). It appears also that  $E_L$  is the lowest possible modulus for a straight line starting at  $A$  and assuring a stiffer material than Eq. (10).

The lowest modulus can be expressed in terms of  $\varepsilon_1$ ,  $\sigma_1$ ; the involved quantities have to satisfy the following relations

$$\varepsilon_1 = \sigma_1/E \quad (11a)$$

$$\varepsilon_2 = (\sigma_2/E) + K(\sigma_2/E)^n \quad (11b)$$

$$E_L = E/[1 + nK(\sigma_2/E)^{n-1}] \quad (11c)$$

$$E_L = (\sigma_2 - \sigma_1)/(\varepsilon_2 - \varepsilon_1) \quad (11d)$$

and the desired solution of Eqs. (11a-d) is obtained, after some algebraic manipulations, as

$$E_L = \frac{E}{1 + NnK(\sigma_1/E)^{n-1}} \quad (12)$$

where

$$N = [n/(n-1)]^{n-1} \quad (13)$$

The lowest modulus  $E_L$  is, of course, smaller than the "tangent modulus" at the stress level  $\sigma_1$

$$E_t = \frac{E}{1 + nK(\sigma_1/E)^{n-1}} \quad (14)$$

and obviously  $E_L < E_t$ . Note also that for ideal elasto-plastic case ( $n \rightarrow \infty$ ) an interesting limit exists

$$\lim_{n \rightarrow \infty} N = e \quad (15)$$

#### Approximate Solution of the Model

Returning to the model of Fig. 1, the exact stress strain relation Eq. (1) is replaced by the two tangents shown in Fig. 4. The analysis is now carried separately for each line.

In segment OA

$$\sigma/E = \theta - \theta_o \quad \theta_o \leq \theta \leq \theta_1 \quad (16)$$

and substituting from Eq. (2)

$$P/P_{cr} = 1 - (\theta_o/\theta) \quad \theta_o \leq \theta \leq \theta_1 \quad (17)$$

Equation (17) describes the load deflection behavior for a linear material.

In segment AC

$$\sigma = E(\theta_1 - \theta_o) + E_L(\theta - \theta_1) \quad \theta \geq \theta_1 \quad (18)$$

but from Eqs. (12) and (16)

$$E_L = \frac{E}{1 + NnK(\theta_1 - \theta_o)^{n-1}} \quad (19)$$

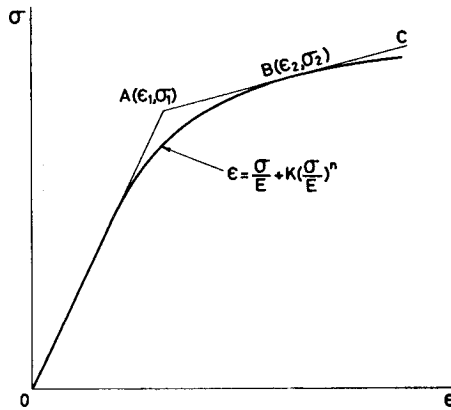


Fig. 4 The lowest modulus.

Combining Eqs. (2, 18, and 19) yields finally

$$\frac{P}{P_{cr}} = \frac{\theta_1 - \theta_o}{\theta} + \frac{1 - (\theta_1/\theta)}{1 + NnK(\theta_1 - \theta_o)^{n-1}} \quad \theta \geq \theta_1 \quad (20)$$

Equations (17) and (20) describe the behavior of a material which is stiffer than the true material. Hence the curve  $P/P_{cr}$  vs  $\theta$  obtained from these equations is an upper bound for the real curve.

Regarding  $\theta_1$  (point A in Fig. 4) as a floating parameter and changing it in small steps, a sequence of curves is obtained. Numerical results for two materials are shown in Figs. 2 and 3. Also, the exact results obtained before are shown.

The envelope formed by the family of the curves given by Eq. (20) can be found in the usual way. Equation (20) is rewritten in the form

$$F(P, \theta, \theta_1) = (P/P_{cr})[1 + NnK(\theta_1 - \theta_o)^{n-1}]\theta - NnK(\theta_1 - \theta_o)^n - \theta + \theta_o = 0 \quad (21)$$

The envelope is obtained<sup>3</sup> by considering, simultaneously, Eq. (21) and the equation

$$\frac{\partial F}{\partial \theta_1} = \frac{P}{P_{cr}} Nn(n-1)K(\theta_1 - \theta_o)^{n-2}\theta - Nn^2K(\theta_1 - \theta_o)^{n-1} = 0 \quad (22)$$

or

$$\theta_1 - \theta_o = [(n-1)/n](P/P_{cr})\theta \quad (23)$$

Substitution of Eq. (23) into Eq. (21) yields the exact solution, Eq. (3). Thus, the envelope of the approximate solutions, Eq. (20), is the true behavior of the model.

The maximum load carrying capacity (for the given model) of the approximated system can be obtained as follows: differentiating Eq. (20) yields

$$\frac{d}{d\theta} \left( \frac{P}{P_{cr}} \right) = \frac{1}{\theta^2} \cdot \frac{\theta_o - NnK(\theta_1 - \theta_o)^n}{1 + NnK(\theta_1 - \theta_o)^{n-1}} \quad (24)$$

Thus, the load increment in the second segment given by Eq. (20) is positive or negative according to the sign of the numerator in Eq. (24). Equating Eq. (24) to zero, the point  $\theta_1 = \theta_1^{**}$ , where there is no load increment from Eq. (20), is obtained

$$\theta_o - NnK(\theta_1^{**} - \theta_o)^n = 0 \quad (25)$$

The solution of Eq. (25) is

$$\theta_1^{**} = \theta_o + [(n-1)/n][\theta_o/(n-1)K]^{1/n} \quad (26)$$

and note that [see Eq. (9)]

$$\theta_1^{**} = [(n-1)/n]\theta^* \quad (27)$$

The maximum load for the approximate system,  $P^{**}$ , is obtained by substituting Eq. (26) into Eq. (20)

$$\frac{P^{**}}{P_{cr}} = 1 / \left[ 1 + nK^{1/n} \left( \frac{\theta_o}{n-1} \right)^{(n-1)/n} \right] \quad (28)$$

which is, of course, the same as the exact solution Eq. (8).

Finally it should be pointed out that the "lowest modulus"  $E_L^{**}$  corresponding to  $P^{**}$ , is equal to the tangent modulus  $E_t^*$

which corresponds to  $P^*$ , and that the following relations hold

$$E_t^*/E = P^*/P_{cr} = P^{**}/P_{cr} = E_L^{**}/E \quad (29)$$

### Conclusions

It has been shown that, for the simple model considered, the envelope of the floating piecewise linear approximation represents the exact nonlinear solution.

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## Measurements in the Near Field of Supersonic Nozzles for Chemical Laser Systems

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### Introduction

AS a contribution to the design of laser systems involving banks of multiple nozzles, the flowfields and gas mixing characteristics near the exits of two Direct Combustion Laser (DCL) nozzle banks were investigated using simulated exhaust gases.<sup>1</sup> The experimental techniques used in the investigation were: 1) laser Doppler velocimeter to measure flow velocities, 2) probes to measure pitot pressure, 3) electron beam to measure both relative and absolute specie number densities, and 4) electron beam to make flow visualization photographs. Also, an analytical investigation was conducted of the flow inside the nozzles, and these results were in good agreement with the experimental data.<sup>1</sup> Only selected results for one nozzle bank are presented in this Note.

### DCL Nozzle Banks

The DCL nozzle banks were supplied by the Air Force Rocket Propulsion Lab. The nozzle bank corresponding to the present data was made of nickel. It had conical nozzles with exit diameters ( $d_e$ ) of 0.234 in., area ratios of 16.85, 10° divergent wall half angles, and 0.024-in.-diam  $H_2$  orifices. A schematic is given in Fig. 1 of the locations of four profiles for which measurements

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